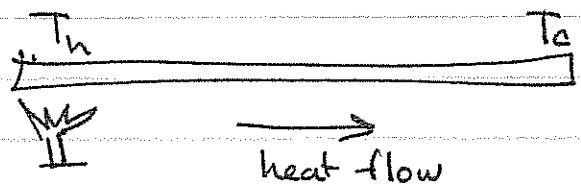


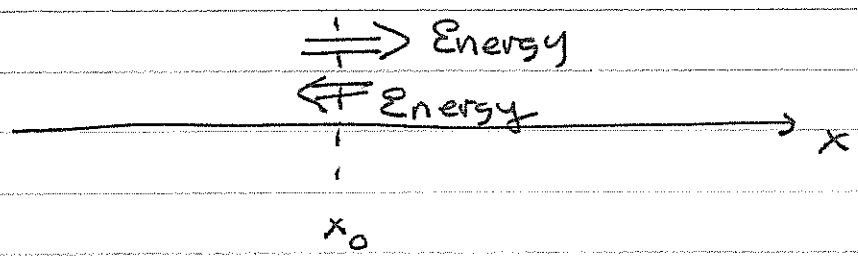
Other Diffusion Processes

① Thermal Conduction:

Temperature gradient drives transfer of energy



Consider flux of energy across x_0 :



$$\text{E-flux} = \left(\frac{\text{no. "particles"}}{\text{area} \cdot \text{time}} \right) \left(\text{average energy per particle} \right)$$

$$= n v \left(\frac{1}{2} k_B T \right)$$

\downarrow $\langle v \rangle$ \uparrow Equipartition (in 1D)

Now $T(x)$, so compare flux at $x_0 \pm l$:

$$\text{Nett E-flux at } x_0 = \frac{1}{2} n v k_B (T(x_0+l) - T(x_0-l))$$

$$\approx \frac{1}{2} n v k_B \left(+2l \frac{dT}{dx} \right)_{x_0}$$

$$\equiv Q \leftarrow \text{heat transferred}$$

Important features:

- the form is $Q = \kappa \frac{dT}{dx}$

where $\kappa =$ thermal conductivity constant

- $\kappa \propto n \langle v \rangle l = \frac{\langle v \rangle}{\sqrt{2} \sigma_0}$

\Rightarrow depends on scattering cross section

\Rightarrow depends on $T^{1/2}$ through $\langle v \rangle$:

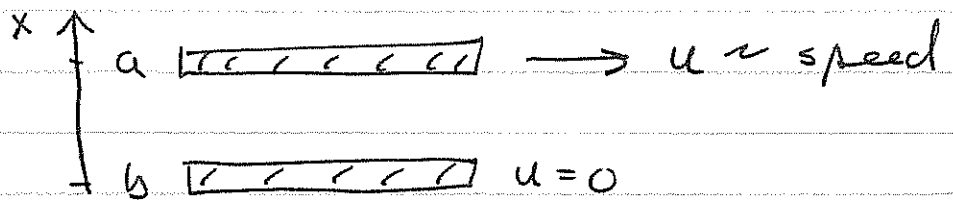
$$\langle v \rangle = \sqrt{\frac{8}{\pi} \frac{k_B T}{m}}$$

for ideal gas of density n_s .

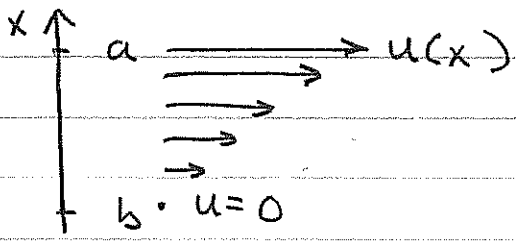
- Note: these 1 dimensional models are artificial. Careful 3 dimensional models produce different numerical pre factors (ie, $\kappa = \frac{1}{3\sqrt{2}} c \langle v \rangle$; $c = (\frac{dE}{dT})_v$).

② Viscosity: diffusion of momentum

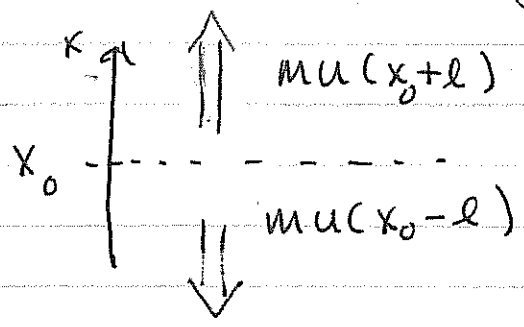
Suppose two plates - one moves



Viscous fluid => top follows plate a, bottom follows plate b, velocity grading:



Fluid has mass, so have momentum gradient and diffusion along x:



Change in momentum => force => pressure
so

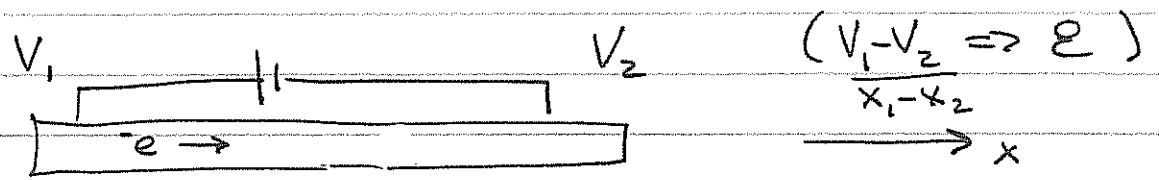
$$P(x_0) = n \langle v \rangle [m u(x_0+l) - m u(x_0-l)]$$

$$\Rightarrow P(x) \Rightarrow \underbrace{n m \langle v \rangle}_{\propto T^{1/2} / \sigma_2} \frac{d}{dx} u(x)$$

=> Experiment confirms that κ & η vary as $T^{1/2}$ — same physics for two different phenomena

③ Electrical Conduction

Current of charge caused by voltage difference:



Here current J along x is

$$J = ne v = ne \langle v \rangle$$

But charges move, between collisions, under E :

$$m \frac{d}{dt} v = eE$$

$$\Rightarrow v = \frac{eE}{m} t + v(0)$$

Now average to form $v = \langle v \rangle$:

$$\langle v \rangle = v = \frac{eE}{m} \langle t \rangle + \langle v(0) \rangle$$

Expect

$$\langle v(0) \rangle = 0 \text{ for random collisions}$$

Know

$$\langle t \rangle = \int_0^{\infty} t \frac{1}{\tau} e^{-t/\tau} dt = \tau$$

So

$$v = \frac{q}{m} E \tau$$

$$\Rightarrow J = \frac{ne^2}{m} E \tau = \sigma_{\text{elec}} E$$

conductivity

or

$$\sigma_{\text{elec}} = \frac{ne^2 \tau}{m} = \frac{ne^2}{m} \frac{1}{\tilde{n} v \sigma_0}$$

no. targets
≠ scatterers

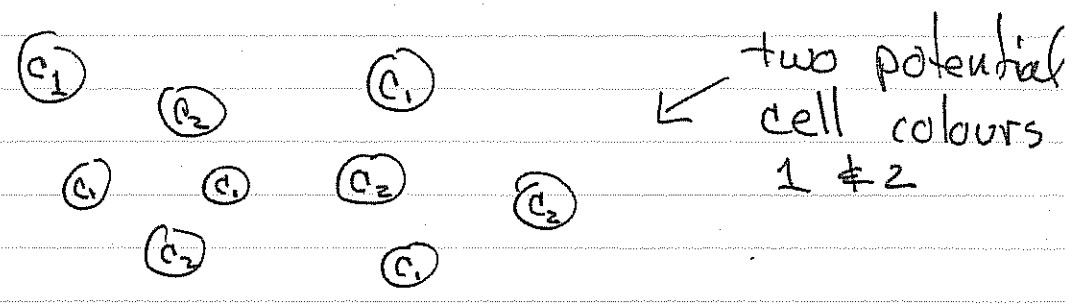
$$\Rightarrow \sigma_{\text{elec}} \propto T^{-1/2}$$

conductivity $\rightarrow \infty$
as $T \rightarrow 0$

④ Pattern formation

Example: skin patterns

One model is that morphogen molecules diffuse across skin with uniform distribution pigment cells (embryo skin)



- Suppose c_1 produce inhibitor morphogens (suppress colour 2 cells) & activators (convert colour 2 cells to colour 1)

- Let inhibitor have larger λ (mean path)

- Net effect on a cell depends on sum of inhibitor & activator concentrations: $n_1 - n_2$

$$\Rightarrow \frac{d}{dt} m(\vec{r}, t) = D \cdot \nabla \cdot \nabla m - km + Q$$

\uparrow morphogen concentration $n_1 - n_2$ \uparrow anisotropic diffusion \uparrow decay } production

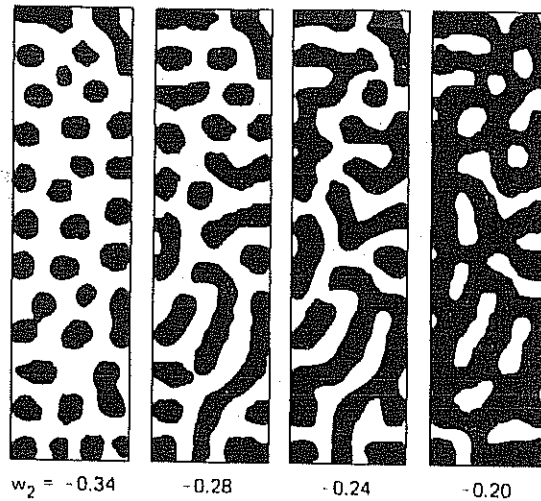


Figure 2.14 Skin Patterns Emerging from Different Inhibitor Values

Example computer simulation of resulting patterns for different thresholds for cell activation.

A General Theory: applicable to dynamic pattern formation (turbulence)

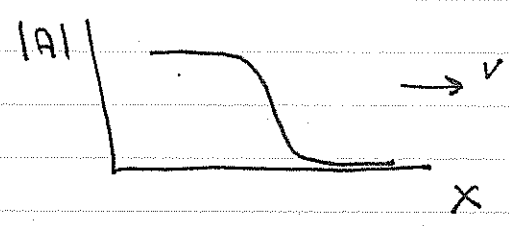
Pattern generation classified by equation type. Example: complex Landau Ginzburg Eq

$$\tau_0 \frac{\partial}{\partial t} A = \epsilon_0^2 \frac{\partial^2}{\partial x^2} A + \epsilon A - g |A|^2 A$$

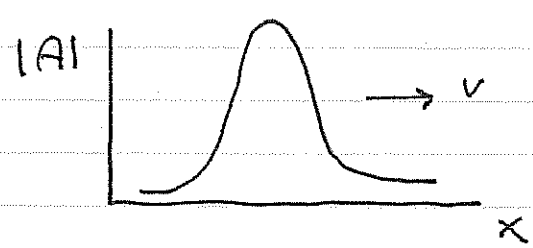
$$A \sim \underbrace{a e^{ik_0 x} + a^* e^{-ik_0 x}}_{\text{Re fields}} + \dots$$

solutions include

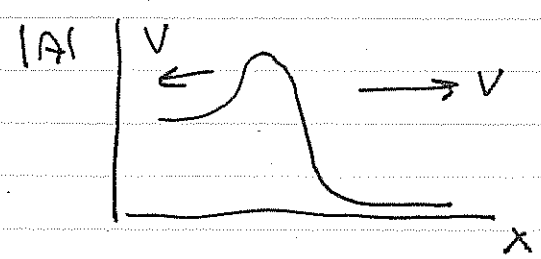
fronts



pulses



sources



sinks

